

STUDY OF REFLECTION AND TRANSMISSION COEFFICIENTS IN SELF –REINFORCED MEDIUM IN THE PRESENCE OF DISSIPATION

#¹**Dr.A.Sinduja**, *Assisstant Professor*,
 #²**Mrs. Ch.Leelavathi**, *Assisstant Professor*,
 #³**Mrs.D.Sridevi**, *Assisstant Professor*,
 #⁴**Mrs. K.Vasavi**, *Assisstant Professor*,

Sai Spurthi Institute of Technology(Autonomous), Sathupalli, Khammam, Telangana

ABSTRACT : The reflection and transmission coefficients of self-reinforced poroelastic medium in the presence of dissipation are calculated. In this the reflection and transmission coefficients of fast dilatational P-I wave , slow dilatational P-II wave and the shear SV wave are calculated. The results are depicted in the graphs which are in numerical session, for this MATLAB is used to get the numerical results.

Keywords: *Reflection Coefficient, Transmission Coefficient, Self-Reinforced Medium, Wave Propagation, Dissipation Effects, Elastic Wave Mechanics.*

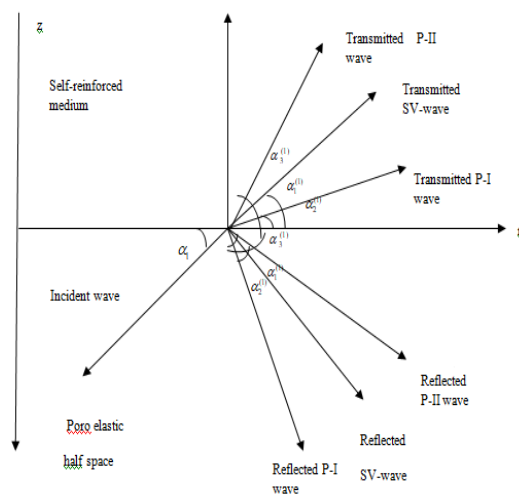
1. INTRODUCTION

During the past half century in many scientific fields, the wave reflection and transmission phenomena have long history and received much attention such as Marine Seismology, Geotechnical Engineering, Acoustics, and Geophysics. Analysis of reflection and transmission phenomena can be used to understand the various materials. Sandstone is a great source for quartz which is very useful in day to day life. Sandstone deposits are may be cylindrical in shape and may be surrounded by poroelastic medium, soil or rock. The study of reflection and transmission of waves incidented at the interface gives the information pertaining to sandstone deposits. The reflection and transmission coefficients are generally based on the medium consisting of two non-homogeneous layers separated by a horizontal interface. In Acoustics, the said coefficients are used to understand the effect of various materials on their acoustic environments. In the research

domains such as Geophysics, and Medicine, the pertinent analysis can be used as Non Destructive Evaluation (NDE) tool. Employing transfer matrix method, Bogy and Gracewski [1] derived the plane wave reflection coefficient for a layered solid half-space, for an isotropic, homogeneous elastic layer. Sinha and Elsibai [2] investigated the reflection of thermoelastic waves from the free surface of a solid half space and at the interface of two semi-infinite media in welded contact. For generalized thermo elastic medium, the reflection phenomena of SV-waves are discussed by Abd-Alla and Al-Dawy [3]. In the paper [3], the reflection coefficient is calculated for shear wave that incident from within the solid on its boundary. Singh [4] exploit the reflection of *P* and *SV* waves from the free surface of an elastic solid with generalized thermo-diffusion. Employing Biot's theory of poroelasticity [5], reflection of plane waves at boundaries of a poroelastic half space is discussed by Tajuddin and Hussaini [6]. In

the paper [6], it is clear that the overlapping parameter between solid and fluid plays a significant role in generating the reflected slow dilatational wave. Dai and Kaung [7] calculated the reflection and transmission of elastic waves at the interface between water and double porosity solids. The analytical expressions of all the three phase velocities of qP, qSV and qSH waves has been derived by Chattopadhyay [8]. Reuven Gordon [9] explained the reflection of cylindrical surface waves. In this the reflection coefficients of fast dilatational ($P-I$) wave, slow dilatational ($P-II$) wave and the shear SV waves with angle of incidence are calculated.

Formulation and Solution of the Problem



Consider a self-reinforced layer, If the wave is incidented in the reinforced layer two dilatational waves and a shear wave are either reflected or transmitted. In this problem reflection only calculated. The equations of motion of a homogeneous, isotropic poroelastic solid in absence of dissipation are [5]

$$N \nabla^2 \vec{u} + \nabla(Ae + Q\varepsilon) = \frac{\partial^2}{\partial t^2} (\rho_{11} \vec{u} + \rho_{12} \vec{U}) + b \frac{\partial}{\partial t} (U - u),$$

$$\nabla(Qe + R\varepsilon) = \frac{\partial^2}{\partial t^2} (\rho_{12} \vec{u} + \rho_{22} \vec{U}) + b \frac{\partial}{\partial t} (U - u)$$

where $\vec{u} = (u, v, w)$ and $\vec{U} = (U, V, W)$ are displacement vectors of solid and fluid, respectively, e and ε are the dilatations of solid and fluid, $P = A + 2N$, Q , and R are all poroelastic constants, b is the dissipation coefficient. The mass coefficients ρ_{11} , ρ_{12} and ρ_{22} satisfy the relations $\rho_1 = \rho_{11} + \rho_{12} = (1 - \beta)\rho_s$,

$$\rho_2 = \rho_{12} + \rho_{22} = \beta\rho_f,$$

$\rho = \rho_1 + \rho_2 = \rho_s + \beta(\rho_f - \rho_s)$ [5]. Here ρ_1 and ρ_2 are mass of solid and fluid per unit volume of aggregate, ρ is total mass of solid-fluid aggregate per unit volume ρ_{11} , ρ_{12} and ρ_{22} are mass coefficients, ρ_s and ρ_f are mass density of the solid and fluid, β is the porosity of the aggregate. The solid stresses σ_{ij} and fluid pressure s are given by

$$\sigma_{ij} = 2Ne_{ij} + (Ae + Q\varepsilon)\delta_{ij}, \quad (i, j = 1, 2, 3), \quad s = Qe + R\varepsilon. \quad (2)$$

In eq. (2), δ_{ij} is the well-known Kronecker delta function.

For an axisymmetric waves, the displacement vectors of solid and fluid are $\vec{u}(u_1, 0, w_1)$ and $\vec{U}(U_1, 0, W_1)$, respectively.

The stress tensor for a self-reinforced layer along preferred direction $\vec{a}$ is given [15]

$$\sigma_{ij} = \lambda e_{kk} \delta_{ij} + 2\mu_T e_{ij} + a(a_k a_l e_{kl} \delta_{ij} + a_i a_j e_{kk}) + 2(\mu_L - \mu_T)(a_i a_k e_{kj} + a_j a_k e_{ki}) + \beta a_k a_m e_{km} a_i a_j \quad (3)$$

here $i, j, k, l, m = 1, 2, 3$, σ_{ij} is the incremental stress component, e_{ij} 's are the strain components given by

$$e_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \text{ and } \vec{a} = (a_1, a_2, a_3)^T$$

is the preferred unit direction of reinforcement. The potential decomposition of displacements for

pertinent upper self reinforced layer as follows:

$$\begin{aligned} u_1 &= \frac{\partial \phi_1}{\partial r} - \frac{\partial \psi_1}{\partial z}, \quad U_1 = \frac{\partial \phi_2}{\partial r} - \frac{\partial \psi_2}{\partial z}, \\ w_1 &= \frac{\partial \phi_1}{\partial z} + \frac{\partial \psi_1}{\partial r} + \frac{\psi_1}{r}, \quad W_2 = \frac{\partial \phi_2}{\partial z} + \frac{\partial \psi_2}{\partial r} + \frac{\psi_2}{r}. \end{aligned} \quad (4)$$

The displacement potentials ϕ 's and ψ 's which are functions of r, z , and t , and are assumed as. [6]

$$\begin{aligned} \phi_1 &= \begin{cases} \phi(1) + \phi_1^{(k)} = A_{1k} e^{i(\omega t - \beta_1 (r \sin \alpha - z \cos \alpha))} + A_1^{(k)} e^{i(\omega t - \beta_1 (r \sin \alpha - z \cos \alpha))}, & l, k=1, 2, \\ \phi(2) + \phi_1^{(k)} = A_{1k} e^{i(\omega t - \beta_1 (r \cos \alpha - z \sin \alpha))} + A_1^{(k)} e^{i(\omega t - \beta_1 (r \cos \alpha - z \sin \alpha))}, & l, k=3, \end{cases} \\ \phi_2 &= \begin{cases} \phi(k) + \phi_2^{(k)} = A_{2k} e^{i(\omega t - \beta_2 \mu_k (r \sin \alpha - z \cos \alpha))} + A_2^{(k)} e^{i(\omega t - \beta_2 \mu_k (r \sin \alpha - z \cos \alpha))}, & l, k=1, 2, \\ \phi(2) + \phi_2^{(k)} = A_{2k} e^{i(\omega t - \beta_2 \mu_k (r \cos \alpha - z \sin \alpha))} + A_2^{(k)} e^{i(\omega t - \beta_2 \mu_k (r \cos \alpha - z \sin \alpha))}, & l, k=3, \end{cases} \\ \omega_1 &= A_3^{(k)} e^{i(\omega t - \beta_3 (r \sin \alpha_1 - z \cos \alpha_1))}, \quad \omega_2 = A_4^{(k)} e^{i(\omega t - \beta_4 (r \sin \alpha_1 - z \cos \alpha_1))}. \end{aligned} \quad (5)$$

Using eq. (4) in eq. (5), after a long calculation, the solid displacements are obtained as follows

$$\begin{aligned}
 u_1 &= \begin{cases} -i\{\delta_{k1}\theta(k)\sin\alpha_k + \sum_{l=1}^2 \delta_{l1}\theta_l^{(k)}\sin\alpha_l^{(k)} - \delta_4\varphi_4^{(k)}\cos\alpha_4^{(k)}\} & k=1,2, \\ i\{\delta_{k1}\theta(k)\cos\alpha_k + \sum_{l=1}^2 \delta_{l1}\theta_l^{(k)}\sin\alpha_l^{(k)} - \delta_4\varphi_4^{(k)}\cos\alpha_4^{(k)}\} & k=3, \end{cases} \\
 w_1 &= \begin{cases} -i\{\delta_{k1}\theta(k)\cos\alpha_k + \sum_{l=1}^2 \delta_{l1}\theta_l^{(k)}\cos\alpha_l^{(k)} + (\theta_{l1}^{(k)}/r) + \delta_4\varphi_4^{(k)}\sin\alpha_4^{(k)}\} & k=1,2, \\ i\{\delta_{k1}\theta(k)\sin\alpha_k + \sum_{l=1}^2 \delta_{l1}\theta_l^{(k)}\cos\alpha_l^{(k)} + (\theta_{l1}^{(k)}/r) + \delta_4\varphi_4^{(k)}\sin\alpha_4^{(k)}\} & k=3. \end{cases}
 \end{aligned} \quad (6)$$

Using these displacement components in eq. (2), the non-zero stress components and fluid pressure are calculated and are given by

$$\begin{aligned}
 & \left\{ \alpha_1 \alpha_3 + 2\alpha_1 \alpha_2 (\mu_L - \mu_T) + \beta \alpha_1^2 \alpha_3 \right\} \\
 & \left\{ \left(\delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^2 \sin^2 \alpha_k - \frac{2}{\ell} \delta_{\frac{k}{\ell}} \rho_{\frac{k}{\ell}}^{(\ell)} \sin^2 \alpha_k^{(\ell)} - \delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^{(2)} \cos 2 \alpha_k^{(\ell)} \right) \right. \\
 & \left. + (\mu_T + (\mu_L - \mu_T) \alpha_1^2 + \alpha_3^2 + \beta \alpha_1^2 \alpha_3^2) \right. \\
 & \left. \left\{ \left(\delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^2 \sin 2 \alpha_k - \frac{2}{\ell} \delta_{\frac{k}{\ell}} \rho_{\frac{k}{\ell}}^{(\ell)} \sin 2 \alpha_k^{(\ell)} - \delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^{(2)} \cos 2 \alpha_k^{(\ell)} \right) \right. \right. \\
 & \left. \left. + (\mu_T + (\mu_L - \mu_T) \alpha_1^2 + \alpha_3^2 + \beta \alpha_1^2 \alpha_3^2) \right\} \right. \\
 & \left. \left\{ \left(\delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^2 \sin 2 \alpha_k - \frac{2}{\ell} \delta_{\frac{k}{\ell}} \rho_{\frac{k}{\ell}}^{(\ell)} \sin 2 \alpha_k^{(\ell)} - (\rho_{\frac{k}{\ell}}^{(\ell)})^2 / r + \delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^{(2)} \cos 2 \alpha_k^{(\ell)} + (i/r) \delta_{\frac{k}{\ell}} \rho_{\frac{k}{\ell}}^{(\ell)} \sin \alpha_k^{(\ell)} \right), \right. \right. \\
 & \left. \left. \left(\delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^2 \cos 2 \alpha_k - \frac{2}{\ell} \delta_{\frac{k}{\ell}} \rho_{\frac{k}{\ell}}^{(\ell)} \sin 2 \alpha_k^{(\ell)} - (\rho_{\frac{k}{\ell}}^{(\ell)})^2 / r^2 + \delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^{(2)} \cos 2 \alpha_k^{(\ell)} + (i/r) \delta_{\frac{k}{\ell}} \rho_{\frac{k}{\ell}}^{(\ell)} \sin \alpha_k^{(\ell)} \right) \right\} \right. \\
 & \left. + \left\{ \left(\delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^2 \sin 2 \alpha_k - \frac{2}{\ell} \delta_{\frac{k}{\ell}} \rho_{\frac{k}{\ell}}^{(\ell)} \sin 2 \alpha_k^{(\ell)} - \delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^{(2)} \cos 2 \alpha_k^{(\ell)} \right) \right\} \right. \\
 & \left. \left\{ \left(\delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^2 \cos 2 \alpha_k - \frac{2}{\ell} \delta_{\frac{k}{\ell}} \rho_{\frac{k}{\ell}}^{(\ell)} \sin 2 \alpha_k^{(\ell)} - \delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^{(2)} \cos 2 \alpha_k^{(\ell)} \right) \right\} \right\} \\
 & \left\{ \alpha a_{13} + 2\alpha a_{13} (\mu_L - \mu_T) + \beta \alpha a_{13}^2 \right\} \\
 & \left\{ \left(\delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^2 \sin^2 \alpha_k - \frac{2}{\ell} \delta_{\frac{k}{\ell}} \rho_{\frac{k}{\ell}}^{(\ell)} \sin^2 \alpha_k^{(\ell)} - \delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^{(2)} \cos 2 \alpha_k^{(\ell)} \right), \right\} \left\{ \tilde{k}_r = 1, 2 \right. \\
 & \left. \left\{ \left(\delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^2 \sin 2 \alpha_k - \frac{2}{\ell} \delta_{\frac{k}{\ell}} \rho_{\frac{k}{\ell}}^{(\ell)} \sin 2 \alpha_k^{(\ell)} - \delta_{\frac{k}{\ell}}^2 \rho_{\frac{k}{\ell}}^{(2)} \cos 2 \alpha_k^{(\ell)} \right), \right\} \left\{ \tilde{k}_r = 3. \right. \right.
 \end{aligned}$$

$$\begin{aligned} & \left\{ (\lambda + \alpha(a_1^2 - a_2^2) + \beta a_3^2) \left[\left(\delta_{k\ell}^{(k)} \sin^2 \alpha_k - \frac{2}{\alpha_k} \delta_{\ell}^{(k)} \sin 2\alpha_k - \frac{\delta_4^{(k)}}{\alpha_k^2} \cos 2\alpha_k \right), \right. \right. \\ & \left. \left. \left(\delta_{k\ell}^{(k)} \sin 2\alpha_k - \frac{\delta_4^{(k)}}{\alpha_k} \sin 2\alpha_k - \frac{\delta_4^{(k)}}{\alpha_k^2} \cos 2\alpha_k \right) \right] \right\} + \\ & (\sigma_{-1}) = \left\{ \alpha_1 \alpha_3 + 2\alpha_1 \alpha_2 (\mu_t - \mu_r) + \beta_1 \alpha_3^2 \right\} \left\{ \delta_{13}^{(k)} \sin 2\alpha_k + \frac{2}{\alpha_k} \delta_{\ell}^{(k)} (\sin 2\alpha_k^2 - \cos^2 \alpha_k^2) - \delta_4^{(k)} \sin \alpha_k^2 \right\} + \\ & \left\{ \delta_{13}^{(k)} \cos^2 \alpha_k + \frac{2}{\alpha_k} \delta_{\ell}^{(k)} (\sin 2\alpha_k^2 - \cos^2 \alpha_k^2) + \delta_4^{(k)} \sin \alpha_k^2 \right\} \\ & \left\{ (\lambda + 2\mu_r + 2\alpha_2^2 + 4\alpha_1^2 (\mu_t - \mu_r) + \beta_1 \alpha_3^2) \right. \\ & \left. \left\{ -\delta_{13}^{(k)} \sin 2\alpha_k + \frac{2}{\alpha_k} \delta_{\ell}^{(k)} (\cos^2 \alpha_k^2 + (1/r) \cos \alpha_k^2 + \sin 2\alpha_k^2) + \delta_4^{(k)} \cos \alpha_k^2 \right\} \right\} \\ & \left. \left\{ \delta_{13}^{(k)} \cos^2 \alpha_k + \frac{2}{\alpha_k} \delta_{\ell}^{(k)} (\cos^2 \alpha_k^2 + (1/r) \cos \alpha_k^2 + \sin 2\alpha_k^2) + \delta_4^{(k)} \cos \alpha_k^2 \right\} \right\} \begin{cases} k=1, l=1, 2 \\ k/l=3 \end{cases} \end{aligned}$$

Boundary Conditions

Continuity of stress components

$$\begin{aligned}(\sigma_{rz})_1 &= 0_{\text{at } z=0}, \\ (\sigma_{zz})_1 &= 0_{\text{at } z=0}.\end{aligned}\quad (11)$$

Continuity of displacement components

$$\begin{aligned}(u)_1 &= 0 \text{ at } z = 0, \\ (w)_1 &= 0 \text{ at } z = 0.\end{aligned}$$

Using these boundary conditions, after some mathematical calculations the ratio of reflected wave to the incident wave is obtained. in this way we get the reflection coefficients of a dilatational waves and a shear wave.

The reflection coefficients are computed by introducing the non-dimensional parameters given below:

$$a_1 = \frac{P}{H}, a_2 = \frac{Q}{H}, a_3 = \frac{R}{H}, a_4 = \frac{N}{H}, d_1 = \frac{\rho_{11}}{\rho}, d_2 = \frac{\rho_{12}}{\rho}, d_3 = \frac{\rho_{22}}{\rho},$$

$$H = P + 2Q + R, \rho = \rho_{11} + 2\rho_{12} + \rho_{22}, \mu = \frac{(\alpha_1 a_1 - \alpha_2 a_2) - (\alpha_1 a_3 - \alpha_2^2) V_k^2}{(\alpha_2 a_3 - \alpha_3 d_2)},$$

$$\mu_k = \frac{(\alpha_1 a_1 - \alpha_2 a_2) - (\alpha_2 a_3 - \alpha_3^2) V_k^2}{(\alpha_2 a_3 - \alpha_3 d_2)}, V_k = \left(\frac{V_0}{V_k} \right)^2, V_l = \left(\frac{V_0}{V_l} \right)^2, k, l, i = 1, 2, 3,$$

are the normalized phase velocities of k^{th} incident and l^{th} reflected waves respectively. In the above $V_0 = \sqrt{\frac{H}{\rho}}$ is reference velocity, V_1 and V_2 are the velocities of dilatational waves of first and second kind respectively, and V_3 is the velocity of shear wave. For illustration

purpose, two types of poroelastic solid cylinders are used, namely, cylinder-I made up of sandstone saturated with water [14], cylinder-II made up of sandstone saturated with kerosene [15]. The physical parameter values are given in Table-1. Employing these values in eq. (9), reflection coefficients are computed as a function of angle of incidence and the results are depicted in figures 2 to 10. The notations R_{cyl-I}, R_{cyl-II}, used in the figures representing the reflection coefficients of cylinder-I, reflection coefficient of cylinder-II, Table-1: Material Parameters

Material Parameters	a_1	a_2	a_3	a_4	d_1	d_2	d_3
Cylinder-I	0.843	0.065	0.027	0.234	0.901	-0.001	0.1
Cylinder-II	0.96	0.006	0.029	0.412	0.876	0	0.124

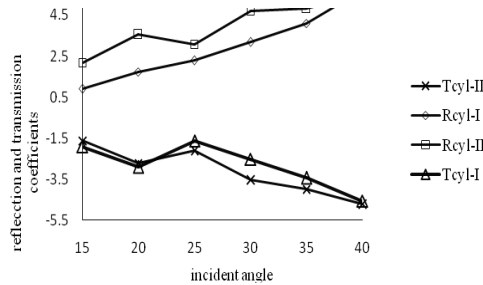


Fig: 2 Reflection and transmission coefficients of $z_1^{(1)}$ and $T_1^{(1)}$.

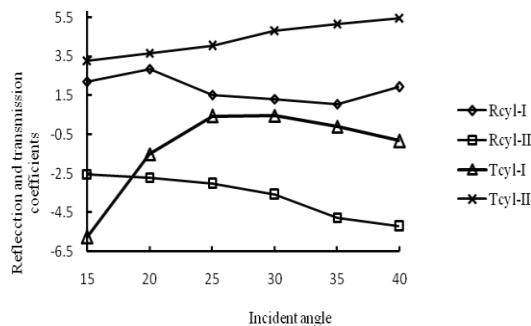


Fig: 3 Reflection and transmission coefficients of $z_1^{(2)}$ and $T_1^{(2)}$.

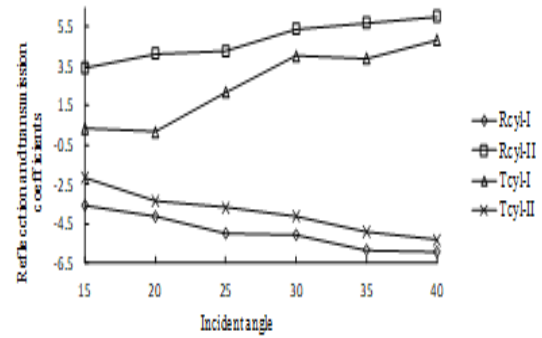


Fig: 4 Reflection and transmission coefficients of $z_1^{(3)}$ and $T_1^{(3)}$.

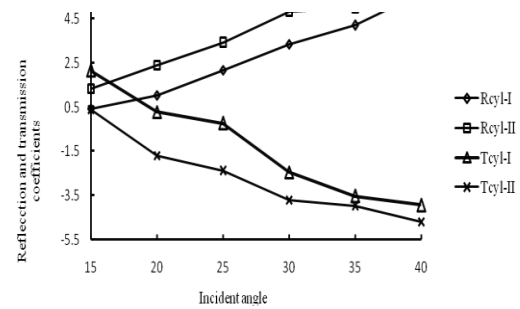


Fig: 5 Reflection and transmission coefficients of $z_2^{(1)}$ and $T_2^{(1)}$.

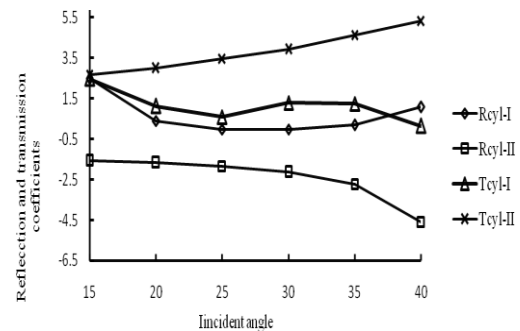


Fig: 6 Reflection and transmission coefficients of $z_2^{(2)}$ and $T_2^{(2)}$.

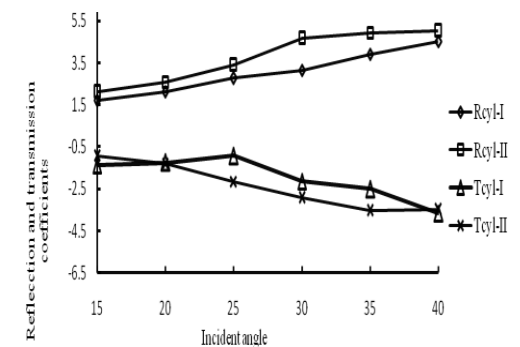


Fig: 7 Reflection and transmission coefficients of $z_2^{(3)}$ and $T_2^{(3)}$.

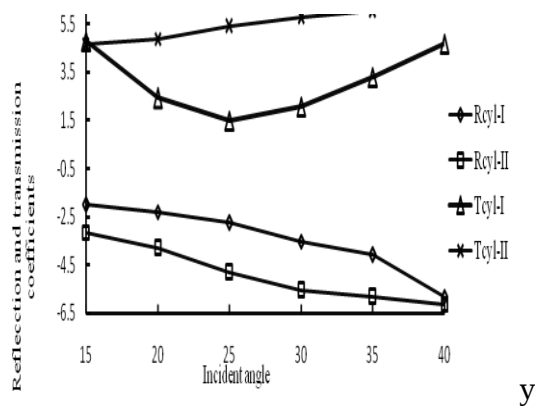


Fig: 8 Reflection and transmission coefficients of $z_3^{(1)}$ and $T_3^{(1)}$.

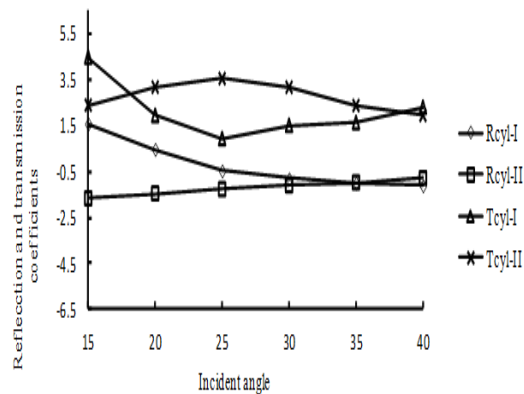


Fig: 9 Reflection and transmission coefficients of $z_3^{(3)}$ and $T_3^{(3)}$.

Incident fast dilatational wave

The reflection coefficient $z_1^{(1)}$ and transmission coefficient $T_1^{(1)}$ of a fast dilatational wave (P-I wave) are calculated. In this the reflection coefficients for cylinder-I and cylinder-II increases gradually as angle of incidence increases.

Incident slow dilatational wave

The reflection coefficient for cylinder-I are periodic and cylinder-II decreases. The reflection coefficient for cylinder-I decreases and cylinder-II increases.

Incident shear SVwave

The reflection coefficients for cylinder -I and cylinder-II increases. in this the reflection coefficient for Cylinder-I decreases, for cylinder-II increases. The transmission coefficient for cylinder-I decreases and for cylinder-II increases.

REFERENCES

- [1] Bogy DB, Gracewski SM (1983) Reflection coefficient for plane waves in a fluid incident on a layered elastic half – space. Journal of Applied Mechanics 50(2): 405-414.
- [2] Sinha SB, Elsibai KA (1996) Reflection of thermoelastic waves at a solid half space with two thermal relaxation times. Journal of Thermal Stress.19:163-777.
- [3] Abd-Alla AN and Al-Dawy AS (2000) Reflection phenomena of SV –waves in generalized thermo elastic medium. International Journal of Mathematics and Mathematical Sciences 23: 529-546.
- [4] Baljeet Singh (2005) Reflection of P and SV waves from the free surface of an elastic solid with generalized thermo diffusion. Journal of Earth system Science 114: 159-168.
- [5] Biot M A (1956) Theory of elastic wave in a fluid-saturated porous solid-I. J Acoust. Soc. America 28: 168-178
- [6] Tajuddin M, Hussaini SJ (2006) Reflection of plane waves at boundaries of a liquid filled poroelastic half space. Journal of Applied Geophysics 58: 59-86.
- [7] Dai ZJ and Kuang ZB (2008) Reflection and transmission of elastic waves at the interface between water and a double porosity solid. Transport in Porous Media 72: 369-392.
- [8] A Chattopadhyay, S Gupta, V.K Sharma, Kumari (2009) Reflection and refraction of plane quasi P waves at a

corrugated interface between distinct triclinic elastic half spaces. International Journal of Solids and Structures. 46 : 3241-3256.

[9] Reuven Gordon (2009) Reflection of Cylindrical surface waves. OPTICS EXPRESS. 17.

[10] Kumar R, Miglani A and Kumar S (2011) Reflection and transmission if plane waves between two different fluid saturated porous half spaces. Bulletin of the Polish Academy of Sciences Technical Sciences 59(2): 227-234.

[11] Mohapatra S, Bora S.N (2011) Reflection and transmission of water waves in a two layer fluid flowing through a channel with undulating bed. Journal of Applied Mathematics and Mechanics. <https://doi.org/10.1002/Zamm.2008>.

[12] Kumar, Sanjeev A, Garg SK (2014) Reflection and transmission of plane waves at the loosely bounded interface of an elastic solid half space and a micro stretch thermo elastic diffusion solid half space. Material Physics and Mechanics 21: 147-167.

[13] Ethyl K, Duwadare QJ, Oyewumi KJ (2016) Semi relativistic reflection and transmission coefficients for two spinless particles. Communication Theory of Physics 66: 389-395.

[14] Malla P, Sindhuja A, Rajitha G (2019) Study of reflection and transmission of axially symmetric body waves incident on a base of a semi-infinite poroelastic solid cylinder. Archive of Applied Mechanics.89.10.1007/S00419-019-01590-5.

[15] Sindhuja A, Rajitha G, MallaReddy P (2020) Shear wave propagation in magneto poroelastic medium sandwiched between Self-reinforced poroelastic medium and poroelastic half-space. Engineering Computations 37(9): 33345-3359.

[16]. Yew CH, Jogi PN (1976) Study of wave motions in fluid-saturated porous rocks. Journal of the Acoustical Society of America 60: 2-8.

[17] Fatt I (1957) The Biot-Willis elastic coefficient for a sand stone. Journal of Applied Mechanics 296-297.